

# Tritium-Concentration Requirements in the Fueling Lines for High- $Q$ Operation in ITER

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One of the most fundamental control problems arising in ITER and future burning-plasma tokamaks is the regulation of the plasma temperature and density to produce a determined amount of fusion power while avoiding undesired transients. ITER is designed to achieve a ratio of fusion power to auxiliary power,  $Q$ , of 10. The reactor's high plasma density ( $> 1 \times 10^{20} m^{-3}$ ) and low burn fraction ( $\sim 1\%$ ) will require high deuterium (D) and tritium (T) fueling rates. Gas puffing will have a low fueling efficiency ( $< 1\%$ ) due to poor neutral penetration [1]. Therefore, the initial phase of ITER will rely on two pellet injectors located on its magnetic high-field-side (HFS) for deep core fueling. The D pellet injector (with pellets of 100% D nominal concentration) and the D-T pellet injector (with pellets of 10%D-90%T nominal concentration) are planned to have maximum throughputs of  $120 Pa m^3/s$  and  $111 Pa m^3/s$ , respectively [2]. However, limitations in the tritium plant subsystems could result in a lower T concentration. Even if the nominal concentration could be initially achieved, it might not be possible to sustain it for the total duration of long pulses. This not only imposes burn-control challenges [3] but also raises serious concerns over having sufficient concentration of T in the fueling lines to sustain long-pulse high- $Q$  operation. In this work, a volume-averaged model of ITER's burning plasma is used to assess the feasibility of accessing  $Q = 10$  operation for different levels of T concentration. Operation points characterized by  $Q = 10$  are sought within ITER's limits for fueling rates and auxiliary heating powers. The results are presented in the form of Plasma Operation CONtour (POPCON) plots that span the density-temperature space. The minimum tritium concentration that can maintain the plasma at  $Q = 10$  is determined for different particle-recycling assumptions. Although this work is concerned with the design parameters of ITER, the analysis can be extended to other future burning-plasma reactors such as DEMO.

## References

- [1] S. K. Combs, L. R. Baylor and others, "Overview of recent developments in pellet injection for ITER," *Fusion Engineering and Design*, vol. 87, pp. 634-640, 2012.
- [2] J.A. Snipes, D. Beltran and others, "Actuator and diagnostic requirements of the ITER Plasma Control System," *Fusion Engineering and Design*, vol. 87, no. 12, pp. 1900-1906, 2012.
- [3] A. Pajares and E. Schuster, "Robust Burn Control in ITER Under Deuterium-Tritium Concentration Variations in the Fueling Lines," *27th IAEA Fusion Energy Conference*, Gandhinagar, India, October 22-27, 2018.

# Tritium-Concentration Requirements in the Fueling Lines for High-Q Operation in ITER

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## Introduction

Tight regulation of the plasma temperature and density to produce a determined amount of fusion power while avoiding undesired transients is required for ITER to operate at  $Q = 10$ , where  $Q$  denotes the fusion power to external power ratio. To attain this goal, the ITER baseline will rely on three heating and current drive (H&CD) systems to supply external power and one pellet injection system (PIS) to provide deuterium (D) and tritium (T) refueling. The three H&CD systems provide 20 MW of electron cyclotron heating, 20 MW of ion cyclotron heating, and 33 MW of neutral beam injection [1]. The PIS consists of two pellet injectors that can fire 5 mm pellets ( $\sim 6 \times 10^{20}$  atoms) at a maximum repetition rate of 16 Hz. The D injector and D-T injector supply, respectively, pellets that have a nominal concentration of 100% D and 10%D-90%T [2]. However, the concentration of T in the fueling lines may fall below its nominal value, particularly during long high- $Q$  pulses. Along with the resulting burn-control challenges [3], this presents concerns over having adequate T refueling to maintain operation at desired states. Accounting for ITER's actuator limits and the T concentration of the pellets, the accessibility of steady-state,  $Q = 10$  operation points is analyzed in this work. The results are determined over the density-temperature space and processed into Plasma Operation CONtour (POPCON) plots. Plasmas with different particle recycling and confinement conditions are compared.

## The Plasma Model

The plasma is modeled by one-dimensional ion density (deuterium density  $n_D$ , tritium density  $n_T$  and alpha particle density  $n_\alpha$ ) and energy density (plasma temperature  $T$ ) rate equations. The radial profiles of the particle densities are assumed to be flat. The radial profile of the temperature has the shape  $T(\vec{r}, t) = T_0(t)[1 - (\phi/\phi_0)]^\gamma$ , where  $T_0$  is the central plasma temperature,  $\phi$  is the toroidal magnetic flux coordinate (at the plasma edge  $\phi = \phi_0$ ) and  $\gamma = 1.85$  [4]. This assumed profile shape approximates the predicted profiles for ITER [5]. The volume average of the one-dimensional equations reduce them to a zero-dimension model given by

$$\frac{d}{dt}n_D = -\frac{n_D}{\tau_D} - n_D n_T \langle \sigma \mathbf{v} \rangle_v + f_{eff} S_D^R + S_D, \quad (1)$$

$$\frac{d}{dt}n_T = -\frac{n_T}{\tau_T} - n_D n_T \langle \sigma \mathbf{v} \rangle_v + f_{eff} S_T^R + S_T, \quad (2)$$

$$\frac{d}{dt}n_\alpha = -\frac{n_\alpha}{\tau_\alpha} + n_D n_T \langle \sigma v \rangle_\nu, \quad (3)$$

$$\begin{aligned} \frac{d}{dt} \left[ \frac{3}{2} nT \right] &= -\frac{3nT_0}{2(1+\gamma)\tau_E} + Q_\alpha n_D n_T \langle \sigma v \rangle_\nu + P_{aux} \\ &+ A_h Z_{eff} I_p^2 a^{-4} T_0^{-3/2} \left( 1 + \frac{3\gamma}{2} \right) - A_b Z_{eff} n_e^2 \frac{T_0^{1/2}}{1+\gamma/2}, \end{aligned} \quad (4)$$

where  $P_{aux}$  is the external heating power,  $S_D$  and  $S_T$  are, respectively, the external fueling rates of deuterium and tritium, and  $Q_\alpha = 3.52$  MeV is the energy deposited into the plasma from each fusion alpha particle. With  $T_0$  expressed in keV, the ohmic heating power ( $P_{ohm}$ ) and bremsstrahlung radiation power ( $P_{brem}$ ) coefficients are, respectively,  $A_h = 2.8 \times 10^{-15}$  MW/m<sup>3</sup> and  $A_b = 5.5 \times 10^{-43}$  MW/m<sup>3</sup> [6].

In addition to the D, T and  $\alpha$  ion species, the plasma is assumed to contain an impurity density  $n_I$  with atomic number  $Z_I$ . The quasi-neutrality condition,  $n_e = n_D + n_T + 2n_\alpha + Z_I n_I$ , defines the electron density. The effective atomic number and the total plasma density are  $Z_{eff} = (n_D + n_T + 4n_\alpha + Z_I^2 n_I)/n_e$  and  $n = n_D + n_T + n_\alpha + n_I + n_e$ , respectively. The tritium fraction of the plasma is defined as  $\gamma = n_T/(n_D + n_T)$ . The fusion reactivity,  $\langle \sigma v \rangle$ , is determined by evaluating at  $T_0$  the formulation found in [7]. To account for the radial temperature profile,  $\langle \sigma v \rangle$  is multiplied by a correction factor,  $G$ , [4] to obtain its volume average,  $\langle \sigma v \rangle_\nu$ .

The IPB98(y,2) H-mode scaling law for the energy confinement time,  $\tau_E$ , is,

$$\tau_E = 0.0562 H_I^{0.93} R^{1.97} B^{0.15} M^{0.19} \epsilon^{0.58} \kappa^{0.78} n_e^{0.41} P_{total}^{-0.69}, \quad (5)$$

where  $H_H$  is the enhancement factor,  $P_{total} = (Q_\alpha S_\alpha + P_{aux} - P_{brem} + P_{ohm}) \times V$  is the total plasma power in MW, and  $V$  is the plasma volume. The ITER values of the parameters in (5) are listed in [8]. The particle confinement times are  $\tau_\alpha = k_\alpha \tau_E$ ,  $\tau_D = k_D \tau_E$  and  $\tau_T = k_T \tau_E$ .

The recycling flux of deuterium and tritium from the reactor walls are modeled as,

$$\begin{aligned} S_D^R &= \frac{1}{1 - f_{ref}(1 - f_{eff})} \left\{ f_{ref} \frac{n_D}{\tau_D} + (1 - \gamma^{PFC}) \times \left[ \frac{(1 - f_{ref}(1 - f_{eff})) R^{eff}}{1 - R^{eff}(1 - f_{eff})} - f_{ref} \right] \left( \frac{n_D}{\tau_D} + \frac{n_T}{\tau_T} \right) \right\}, \\ S_T^R &= \frac{1}{1 - f_{ref}(1 - f_{eff})} \left\{ f_{ref} \frac{n_T}{\tau_T} + \gamma^{PFC} \times \left[ \frac{(1 - f_{ref}(1 - f_{eff})) R^{eff}}{1 - R^{eff}(1 - f_{eff})} - f_{ref} \right] \left( \frac{n_D}{\tau_D} + \frac{n_T}{\tau_T} \right) \right\}, \end{aligned}$$

where  $f_{eff}$ ,  $f_{ref}$ ,  $R^{eff}$  and  $\gamma^{PFC}$  are constants that define the quality of the recycling [9].

The L-H mode transition threshold power, expressed in MW, for a D-T plasma is given by  $P_{thresh} = 4.3 M^{-1} B^{0.772} n_e^{0.782} R^{0.999} a^{0.975}$ . The total plasma power must be above this threshold power for the plasma to be in H-mode [10]. The ratio of the fusion power to the power externally put into the plasma is  $Q = E_{fus} n_D n_T \langle \sigma v \rangle_\nu / (P_{aux} + P_{ohm})$  where  $E_{fus} = 17.6$  MeV.

The D-T pellet injector fuels the plasma at a rate of  $S_{DT}^{line}$  with pellets containing a tritium concentration of  $\gamma_{DT}^{line}$ . The D pellet injector fuels the plasma at a rate of  $S_D^{line}$  with pellets containing a tritium concentration of  $\gamma_D^{line}$ . Therefore, the external fueling rates,  $S_D$  and  $S_T$ , are formulated as,

$$S_D = (1 - \gamma_{DT}^{line})S_{DT}^{line} + (1 - \gamma_D^{line})S_D^{line}, \quad (6)$$

$$S_T = \gamma_{DT}^{line}S_{DT}^{line} + \gamma_D^{line}S_D^{line}. \quad (7)$$

### POPCON Analysis Results

In the following analysis, the steady-state plasma conditions for  $Q = 10$  operation are studied in the density-temperature space. The system of four equations (1)-(4) in steady-state ( $d/dt = 0$ ) is solved for four unknowns. The unknowns are the alpha particle density fraction  $f_\alpha = n_\alpha/n_e$ , the external fueling rate of deuterium  $S_D$ , the external fueling rate of tritium  $S_T$ , and the impurity density fraction  $f_I = n_I/n_e$ . With  $Q = 10$ , the auxiliary power,  $P_{aux}$ , is determined from the expression for  $Q$ . The solution is found repeatedly over a grid of fixed  $n_e$  and  $T_0$  values. All of the plasma conditions (i.e.  $\tau_E$ ) can be calculated from the four unknowns, the electron density and the central temperature. The resulting contour lines are plotted in the  $n_e$ - $T_0$  space.

The two scenarios discussed in this section consider only  $Q = 10$  operation and the following fixed constants:  $\gamma_D^{line} = 0$ ,  $H_H = 1$ ,  $\gamma = 0.5$ ,  $Z_I = 10$  (Neon),  $f_{ref} = 0.5$ ,  $R_{eff} = 0.6$  and  $\gamma^{PFC} = 0.5$ . Fig. 1 shows the contours of constant auxiliary power for a plasma with recycling ( $f_{eff} = 0.2$ ) and high particle confinement ( $k_D = k_T = 1.2$ ,  $k_\alpha = 1.6$ ). Fig. 2 shows the contours of constant fusion power for a plasma with no recycling ( $f_{eff} = 0$ ) and low particle confinement ( $k_D = k_T = k_\alpha = 1$ ). In both figures, multiple lines mark the saturation curves of the ITER actuators. The dashed-blue line marks where  $P_{aux} = 0$ . Below this line, the auxiliary power is negative and the results are physically meaningless. The solid-blue line denotes where  $P_{aux} = 73$  MW. The yellow line marks where the D injector saturates at 16 Hz with  $\gamma_{DT}^{line} = 90\%$ . Where

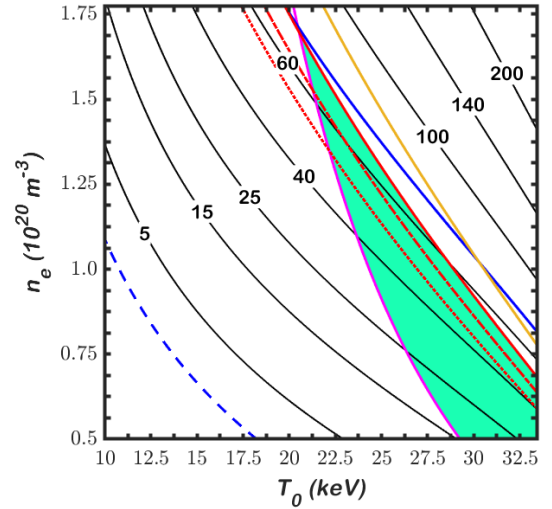


Figure 1: Results with auxiliary power [MW] isolines (black) for plasma with  $k_D = k_T = 1.2$ ,  $k_\alpha = 1.6$  and  $f_{eff} = 0.2$ .

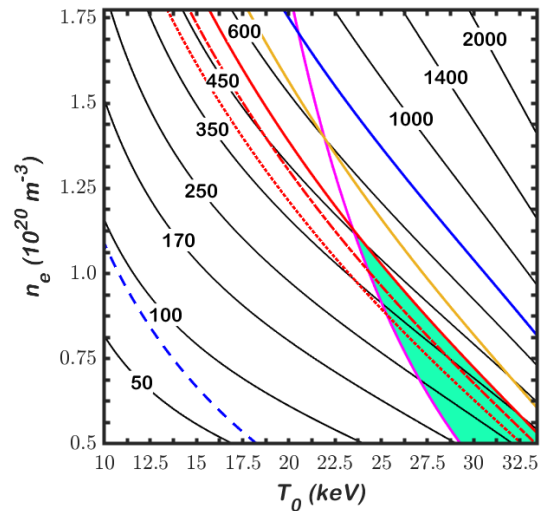


Figure 2: Results with fusion power [MW] isolines (black) for plasma with  $k_D = k_T = k_\alpha = 1$  and  $f_{eff} = 0$ .

the D-T injector hits 16 Hz with  $\gamma_{DT}^{line}=90\%$ ,  $\gamma_{DT}^{line}=80\%$  and  $\gamma_{DT}^{line}=70\%$  is plotted using the solid-red, dashed-dotted-red and dotted-red lines, respectively. On the magenta line, the total plasma power,  $P_{total}$ , equals the threshold power  $P_{thresh}$  (the plasma is in H-mode above this line). The green zone marks the steady-state operable space with  $Q = 10$  for the given ITER plasma ( $\gamma_{DT}^{line} = 90\%$ ). Here, the actuator saturation limits are not violated and the plasma is in H-mode.

### Conclusions and Future Work

Due to enhanced confinement and recycling conditions, the plasma plotted in Fig. 1 has a much larger  $Q = 10$  operable space than that plotted in Fig. 2. In both figures, the three red lines indicating D-T injector saturation (16 Hz) show how a decreasing T concentration ( $\gamma_{DT}^{line}$ ) can shrink the  $Q = 10$  operable space. The steady-state,  $Q = 10$  operable spaces of the ITER plasmas depicted in Fig. 1 and Fig. 2 vanish when  $\gamma_{DT}^{line}$  drops below 30% and 60%, respectively. Future work may include extending this study to other  $Q$  conditions, studying sensitivity of the results, and applying this analysis to other burning reactors such as DEMO.

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### References

- [1] J.A. Snipes, D. Beltran, et al., "Actuator and diagnostic requirements of the ITER Plasma Control System," *Fusion Engineering and Design*, vol. 87, no. 12, pp. 1900-1906, 2012.
- [2] S. K. Combs, L. R. Baylor, et al., "Overview of recent developments in pellet injection for ITER," *Fusion Engineering and Design*, vol. 87, pp. 634-640, 2012.
- [3] A. Pajares and E. Schuster, "Robust Burn Control in ITER Under Deuterium-Tritium Concentration Variations in the Fueling Lines," *27th IAEA Fusion Energy Conference*, Gandhinagar, India, October 22-27, 2018.
- [4] J. J. Martinell and J. E. Vitela, "An Optimal Burn Regime in a Controlled Tokamak Fusion Power Plant," *IEEE Transactions on Plasma Science*, vol. 44, no. 3, pp. 296-305, 2016.
- [5] M. Shimada, A. Chudnovskii, et al., "Physics Design of ITER-FEAT," *J. Plasma Fusion Res.*, vol. 3, pp. 77-83, 2000.
- [6] W. M. Stacey, "Fusion: An introduction to the physics and technology of magnetic confinement fusion," Wiley-VHC, Weinheim, 2nd edition, 2010.
- [7] H. S. Bosch and G. M. Hale, "Improved formulas for fusion cross-sections and thermal reactivities," *Nucl. Fusion* vol. 32 pp. 611, 1992.
- [8] International Atomic Energy Agency, "Summary of the ITER Final Design Report," *ITER EDA Documentation Series*, no. 22, Vienna, 2001.
- [9] J. Ehrenberg, "Wall effects on particle recycling in tokamaks," in *Physical Processes of the Interaction of Fusion Plasmas with Solids*, p. 35, Academic Press, 1996.
- [10] Y. R. Martin, et al., "Power requirement for accessing the H-mode in ITER," *J. Phys.: Conf. Ser.*, vol. 123, no. 1, 2008.